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**Solar origin of heliospheric magnetic field inversions:  
Evidence for coronal loop opening within  
pseudostreamers**

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9 **Abstract.** The orientation of the heliospheric magnetic field (HMF) in  
10 near-Earth space is generally a good indicator of the polarity of HMF foot  
11 points at the photosphere. There are times, however, when the HMF folds  
12 back on itself (is inverted), as indicated by suprathermal electrons moving  
13 sunward while carrying the heat flux away from the Sun. Analysis of the near-  
14 Earth solar wind during the period 1998-2011 reveals that inverted HMF is  
15 present approximately . Inverted HMF is mapped to the coronal source sur-  
16 face, where a new method is used to estimate coronal structure from the potential-  
17 field source-surface model. We find a strong association with bipolar stream-  
18 ers containing the heliospheric current sheet, as expected, but also with unipo-  
19 lar or pseudostreamers, which contain no current sheet. Because large-scale  
20 inverted HMF is a widely-accepted signature of interchange reconnection at  
21 the Sun, this finding provides strong evidence for models of the slow solar  
22 wind which involve coronal loop opening by reconnection within pseudostreamer  
23 belts as well as the bipolar streamer belt. Occurrence rates of bipolar- and  
24 pseudostreamers suggest that they are equally likely to result in inverted HMF  
25 and, therefore, presumably undergo interchange reconnection at approximately  
26 the same rate. Given the different magnetic topologies involved, this suggests  
27 the rate of reconnection is set externally, possibly by the differential rota-  
28 tion rate which governs the circulation of open solar flux.

## 1. Introduction

large-scale heliospheric magnetic field (HMF) is generally well described by the Parker spiral.  $135^\circ/315^\circ$  for outward/inward polarity HMF [e.g., *Borovsky, 2010*]. The heliospheric current sheet (HCS) separates sectors of inward and outward magnetic flux and projects back to a coronal source-surface as a neutral line marking the heliomagnetic equator. Crossings of the near-Earth HCS can be identified by rapid changes in the HMF direction from  $135^\circ$  to  $315^\circ$ , or vice versa. This is shown schematically in Figure 1a.

HMF connectivity to the Sun can usually be inferred by suprathermal electron (STE) observations. Open HMF, which has one end connected to the Sun, exhibits an adiabatically focussed STE beam, or "strahl," that originates in the solar corona [*Feldman et al., 1975; Rosenbauer et al., 1977*]. Thus outward (inward) magnetic sectors should contain a strahl which is parallel (antiparallel) to the HMF, as shown in Figure 1a. , when both parallel and antiparallel strahls are present, reveal "closed" HMF, with both ends of the field line connected to the Sun (times 2 and 3 in Figure 1b). They are strongly associated with interplanetary coronal mass ejections [*Gosling et al., 1987; Wimmer-Schweingruber et al., 2006*], which in turn are frequently encountered at magnetic sector boundaries [*Crooker et al., 1998*, see also Figure 1b].

There also exist periods with a single strahl in the opposite sense to that expected from the magnetic field direction [*Kahler and Lin, 1994, 1995; Kahler et al., 1996; Crooker et al., 1996; Crooker et al., 2004b*], as shown in Figures 1c-e. These intervals imply that the magnetic field is folded back upon itself, or inverted. Inverted HMF intervals can be bounded by a change in the magnetic field direction with no change in the strahl direction

and vice versa. Pairs of the former are common and can be found both near the HCS, as in Figure 1c, and in unipolar regions [e.g., *Balogh et al.*, 1999], as in Figure 1d. These pairs of field changes bound inversions that are usually of short duration, on the order of an hour or two. In contrast, inversions bounded on at least one side by a change in the strahl direction with no change in the magnetic field direction are less common but can be of long duration, on the order of a day or more [*Crooker et al.*, 2004b]. Moreover, they can only be understood in terms of a three-dimensional structure. In cases involving the HCS, as in Figure 1e, where the dashed field lines lie out of the plane of the Figure, the inversion results in a mismatch between the magnetic and electron signatures of the sector boundary [*Crooker et al.*, 2004b].

While some of the smaller inversions may be the product of large-scale turbulent processes, the larger inversions appear to be robust signatures of near-Sun magnetic interchange reconnection, as sketched in Figures 1c-e, where a green X marks a reconnection site. The legs of large loops expanding into the heliosphere reconnect with adjacent open field lines. *Crooker et al.* [2004b] suggest that the expanding loops are at the quiet end of a spectrum of large-scale transient outflows, with coronal mass ejections (CMEs) at the active end. This interpretation is supported by the observation of coronal inflows and collapsing loops at locations where the HCS is inclined to the solar rotation direction [*Sheeley and Wang*, 2001], taken to be signatures of magnetic reconnection. The association of inverted HMF with the HCS suggests the solar origin of the expanding loops can be bipolar helmet streamers which surround the coronal source-surface neutral line and separate magnetic flux from coronal holes of opposite magnetic polarity, e.g., the

two polar coronal holes at solar minimum. This paper also considers unipolar streamers, called "pseudostreamers," as an additional source.

Pseudostreamers are very similar to bipolar streamers in coronagraph observations. They are also formed at the boundary between coronal magnetic flux from two different coronal holes, but unlike bipolar streamers, the flux at both foot points is of the same polarity and, thus, they do not contain current sheets [e.g., *Eselevich*, 1998; *Eselevich et al.*, 1999; *Zhao and Webb*, 2003; *Wang et al.*, 2007]. There has recently been much interest in pseudostreamers as a possible source of the slow solar wind [*Crooker et al.*, 2012; *Riley and Luhmann*, 2012], either through the expansion of coronal magnetic flux tubes [*Wang et al.*, 2012], or through the intermittent release of plasma by the opening of coronal loops via magnetic reconnection [*Antiochos et al.*, 2011]. *Crooker et al.* [2012] demonstrate that pseudostreamers occur in belts which are topologically connected to the bipolar streamer belt, thus forming a network of slow solar wind sources.

In this study we investigate the properties and solar origin of inverted heliospheric magnetic flux during the period 1998 to 2011, for which almost continuous HMF and STE data are available from the Advanced Composition Explorer (ACE) spacecraft. In particular, comparisons are made with the locations of bipolar and pseudostreamers estimated using the potential-field source-surface (PFSS) model of the corona.

## 2. Detection of HMF inversions

The 272eV energy channel is used, as it is well within the suprathermal range, showing little contribution from the core electron population, but still providing high count rates [e.g., *Anderson et al.*, 2012]. The SWEPAM PAD data are available from January 1998 to August 2011, which determines the interval used in this study.

discriminate between closed HMF and  $90^\circ$  pitch-angle depletions owing to mirroring from large-scale, downstream structures [Gosling *et al.*, 2001], so closed flux occurrence is likely overestimated. Furthermore, while counterstreaming electron intervals are separated out from inverted and uninverted flux, no attempt is made to explicitly exclude ICMEs. Indeed, if ICMEs contain "open" inverted field lines, they must result from reconnection in the corona in the same way as ambient solar wind intervals [Owens and Crooker, 2006, 2007]. By including all solar wind data in the study, no assumptions are made about the source and processes involved in the creation of inverted HMF.

There are

### 3. Properties of HMF inversions

Figure 3 shows the probability distribution functions (PDFs) of solar wind parameters. The solar wind properties of

### 4. Association with bipolar and pseudostreamers

Thus to aid in the interpretation of these data, we use a potential-field source-surface (PFSS) model of the corona [Schatten *et al.*, 1969] based on WSO magnetograms to identify the locations of the HCS and, hence, bipolar streamers as well as pseudostreamers.

#### 4.1. Case studies

The pink and light grey regions show, respectively, outward and inward polarity coronal holes, i.e., the photospheric foot points of magnetic field lines reach the source surface at 2.5 solar radii. Red (white) lines show the . Overlaid on the ecliptic plane is the observed magnetic polarity in near-Earth space, ballistically mapped back to the source surface using the observed solar wind speed, with red/white dots indicating , as determined in



113 Section 2. For this particular Carrington rotation, there is agreement between the mag-  
114 netic polarity predicted by the PFSS model and that observed near-Earth. Green crosses  
115 show the coronal source-surface locations of observed HMF inversions at the heliographic  
116 latitude of Earth.

117 The two intervals of inverted HMF at Carrington longitude of . The remaining HMF  
118 inversions are also associated with a change in magnetic connectivity, with the photo-  
119 spheric foot points along Earth orbit shifting between different coronal holes, but without  
120 an associated change in foot point polarity, indicative of pseudostreamers. These HMF  
121 inversions are thus associated with pseudostreamers rather than bipolar streamers.

122 define a parameter  $dS$ , the distance between photospheric foot points of neighbouring  
123 points on the source surface. In practice, the magnitude of  $dS$  will depend on the spatial  
124 resolution at which field lines are traced, making units somewhat arbitrary. In this study,  
125 we calculate  $dS$  by moving along the ecliptic plane in  $1^\circ$  steps. When adjacent points on  
126 the source surface map to the same coronal hole,  $dS$  will be small, for example as seen  
127 between  $0^\circ$  and  $60^\circ$  Carrington longitude for CR1990. When neighbouring source-surface  
128 points map to different coronal holes, however, such as the HCS crossing at  $310^\circ$  Carring-  
129 ton longitude,  $dS$  will be very large. The middle panel of Figure 4 shows  $\log_e(dS)$  as a  
130 function of Carrington longitude along the ecliptic plane. Vertical yellow lines mark HCS  
131 crossings, where  $\log_e(dS)$  spikes correspond to bipolar streamers. The dashed horizontal  
132 line at  $\log_e(dS) = 3$  marks the threshold selected to define a streamer. It is the value  
133 which  $\log_e(dS)$  reaches or exceeds at all HCS crossings in the 1998 to 2011 period and  
134 corresponds to source surface points with a  $1^\circ$  separation having a photospheric footpoint  
135 separation of  $\geq 5^\circ$ . It thus selects all bipolar streamers and appears to select most sig-

nificant pseudostreamers while suppressing smaller structures. Blue vertical lines mark  $\log_e(dS)$  spikes without polarity reversals, our definition of a pseudostreamer. The 17 1-hour intervals of inverted HMF not associated with the HCS in CR1990 all map close to the longitudes of pseudostreamers.

The bottom panel of Figure 4 is a contour plot of  $dS$  at all latitudes. It demonstrates in another way the finding reported by [Crooker *et al.*, 2012] that pseudostreamer belts, but connect to the bipolar streamer belt to form a network of slow solar wind sources that expands to cover the source surface during solar maximum. As is the case for bipolar streamers, HMF inversions are not associated with all pseudostreamers; however, Figure 4 demonstrates that streamer-associated inverted HMF is likely to be common at all latitudes near solar maximum.

## 4.2. Statistical analysis

In order to systematically analyse the entire 1998-2011 interval, and define strict thresholds for association between inverted HMF and streamers. We begin by including only Carrington rotations in which the PFSS model provides a reasonable representation of the observed magnetic structure of the corona and solar wind. By assigning +1 (-1) to outward (inward) Parker spiral polarity, and ignoring undetermined, counterstreaming and inverted intervals, we compute the mean-square error (MSE) between the PFSS and observed sector structure mapped to the source surface. Thus MSE is a combination of errors in the PFSS solution and errors in the simple ballistic mapping of near-Earth solar wind to the coronal source surface.

156 of ecliptic longitudes are covered by pseudostreamers (bipolar streamers). Note that the  
157 association scheme allows a single inverted HMF interval to map to both a bipolar and  
158 pseudostreamer if they are located close in longitude. Table 2 summarises these results.

159 In general, there are insufficient inverted HMF events to detect significant differences in  
160 the of solar wind properties of bipolar- and pseudostreamer-associated inversions. Prob-  
161 ability distributions of density, however (not shown), suggest that HMF inversions from  
162 bipolar streamers contain denser solar wind than inverted HMF from pseudostreamers,  
163 consistent with general properties of pseudostreamer-associated solar wind [*Wang et al.*,  
164 2012].

## 5. Conclusions and Discussion

165 The polarity of the photospheric foot point of heliospheric magnetic flux (HMF) can  
166 be independently estimated from both the local HMF orientation, as measured using in  
167 situ magnetometer observations, and the direction of the suprathermal electron beam,  
168 or "strahl." For the bulk of the solar wind, these two methods show agreement. There  
169 are intervals, however, in which the strahl is directed towards the Sun, implying that the  
170 magnetic field line is inverted, or folded back on itself. This is an expected signature of  
171 near-Sun magnetic reconnection by which the Sun can open previously closed heliospheric  
172 loops [*Owens et al.*, 2011; *Owens and Lockwood*, 2012]. Using an automated data analysis  
173 method, we find inverted flux in approximately 5.5% of the solar wind data between 1998  
174 and 2011, though this is likely an underestimate due to strict selection criteria. We do  
175 not find a strong solar cycle variation in the occurrence rate of inverted HMF, but this  
176 finding is confined to the ecliptic plane . Inverted HMF is associated with dense, slow,  
177 cool solar wind, with lower than average magnetic field intensity. In order to determine

178 the solar origin of these structures, we used a potential-field source-surface model to  
179 infer the global structure of the coronal magnetic field and a new automated detection  
180 method for bipolar and pseudostreamers. Of the 2263 1-hour inverted HMF intervals  
181 identified in the solar wind and mapped back to the coronal source surface, 1310 (58%)  
182 are associated with streamers. Given that the probability of a solar wind interval being  
183 associated with a streamer by chance is 52%, the association between inverted HMF  
184 and streamers is significant at the 99.9% level. Of the 1310 streamer-associated inverted  
185 HMF intervals, 949 (504) map to pseudostreamers (bipolar streamers). This ratio is in  
186 reasonable agreement with the occurrence rates of pseudostreamers and bipolar streamers  
187 in the ecliptic plane, 39% and 20%, respectively,

188 If we assume that inverted HMF is primarily a signature of reconnection in the corona  
189 [e.g., *Titov et al.*, 2011], our results suggest that the rate of reconnection is similar within  
190 bipolar and pseudostreamers. This seems reasonable in view of their magnetic structure.  
191 For the bipolar streamer case, a three-dimensional magnetic configuration for interchange  
192 reconnection that can create the inversion is illustrated in 1e and has already been dis-  
193 cussed in section 1. For the pseudostreamer case, an appropriate magnetic configuration  
194 can be drawn in just two dimensions, as illustrated in Figure 6. Closed loops within one  
195 of the two arcades that form pseudostreamers are shown to rise as a result of photospheric  
196 flux emergence, but could equally be the result of loop foot point shearing, etc. In the  
197 top panel, the rising loop undergoes interchange reconnection before it reaches the solar  
198 wind acceleration height and therefore doesn't result in the generation of inverted HMF.  
199 This configuration is common from the solar perspective [e.g., *Wang et al.*, 2007; *Crooker*  
200 *et al.*, 2012]. In contrast, from the heliospheric perspective, the rising loops are dragged

201 out by the solar wind before interchange reconnection takes place, which does generate  
 202 inverted HMF, as illustrated in the bottom panels. Thus pseudostreamer loop expansion  
 203 and opening via interchange reconnection would transport pre-existing open solar flux in  
 204 much the same way as the CME-driven transport proposed by *Owens et al.* [2007]. Indeed,  
 205 as proposed by *Crooker et al.* [2004b] for loops expanding from the helmet arcade in the  
 206 case of bipolar streamers, loops that create inversions from pseudostreamers can also be  
 207 considered as the quiet end of a spectrum of loops, where the active end is CMEs. This  
 208 analogy holds because pseudostreamers are well-documented sources of CMEs [*Fainshtein,*  
 209 1997; *Eselevich et al.*, 1999; *Zhao and Webb*, 2003; *Liu and Hayashi*, 2006].

210 In addition, similar levels of association between inverted HMF with bipolar and pseu-  
 211 dostreamers, despite the differing magnetic topologies, suggest that the reconnection rate  
 212 is externally controlled. One possibility is the stress between the differential rotation of  
 213 the photosphere and the rigid corotation of the corona [*Nash et al.*, 1988; *Wang and Shee-*  
 214 *ley*, 2004] and the consequent circulation of open solar flux [*Fisk et al.*, 1999; *Fisk and*  
 215 *Schwadron*, 2001]. We note that inverted HMF is the expected heliospheric signature of  
 216 large coronal loop opening, one of the proposed mechanisms for slow solar wind formation  
 217 [e.g., *Fisk*, 2003]). Thus our results provide support for the idea of pseudostreamers being  
 218 a source of slow solar wind through intermittent release from previously closed coronal  
 219 loops [*Antiochos et al.*, 2011], though the effect of magnetic flux tube expansion [*Wang*  
 220 *et al.*, 2012] may still be important.

221 Inverted HMF has direct implications for in situ spacecraft estimates of the total mag-  
 222 netic flux threading the solar source surface, often referred to as the unsigned open solar  
 223 flux, OSF [e.g., *Owens et al.*, 2008a]. Figures 1c and 1d clearly illustrate the issue: In-

224 verted HMF provides magnetic flux which threads the heliocentric sphere at 1 AU, but  
 225 does not map back to the source surface, resulting in an overestimate in OSF from in  
 226 situ observations. , decomposing the HMF along the Parker spiral direction, which can  
 227 successfully remove the effects of waves and turbulence [Erdős and Balogh, 2012], may  
 228 not address this particular issue. Both the occurrence rate and magnetic field strength  
 229 associated with inverted HMF are small, suggesting this may not have a large effect on  
 230 OSF estimates. Even if inverted HMF has an average magnetic flux density as high as  
 231 the rest of the solar wind, the decrease in the unsigned OSF would only be  $2 \times 5\% = 10\%$ .  
 232 The factor 2 arises as follows: if inverted HMF intervals contain  $\phi_I$  of magnetic flux, the  
 233 unsigned OSF will be overestimated by  $2\phi_I$ , since both the inverted and "return" flux  
 234 thread the heliocentric surface but not the coronal source surface. We note that, in gen-  
 235 eral, inverted HMF intervals are less than a day long, though this may be partly due to  
 236 the strict criteria used and the time interval considered [c.f. Crooker et al., 2004b]. Thus  
 237 taking 1-day averages of the radial magnetic field for the purposes of estimating OSF  
 238 may indirectly negate the effect of inverted HMF [c.f. Wang and Sheeley, 1995], though  
 239 it does not directly address the issue of physical origin [see also Lockwood et al., 2009, for  
 240 discussion of correction of 1-AU measurements to the coronal source surface].

241 In summary, we have developed a new method for identifying bipolar streamers and  
 242 pseudostreamers in PFSS synoptic maps. The results confirm that together these struc-  
 243 tures form a network of slow solar wind sources which expands over the source surface at  
 244 solar maximum. Moreover, we have analyzed suprathermal electron data from the solar  
 245 wind and find that, like bipolar streamers, pseudostreamers are sources of HMF inversions.  
 246 These are understood to be signatures of coronal loops that expand into the heliosphere

and subsequently become open through reconnection in the corona. Loop-opening is a key process in one of two competing models for the source of the slow wind.

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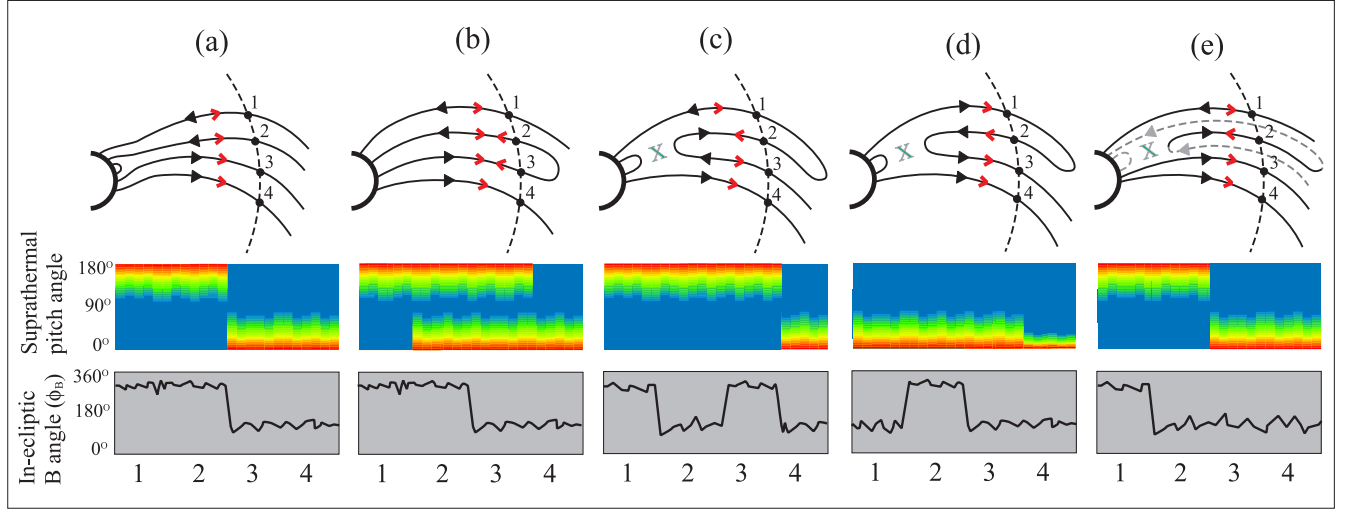
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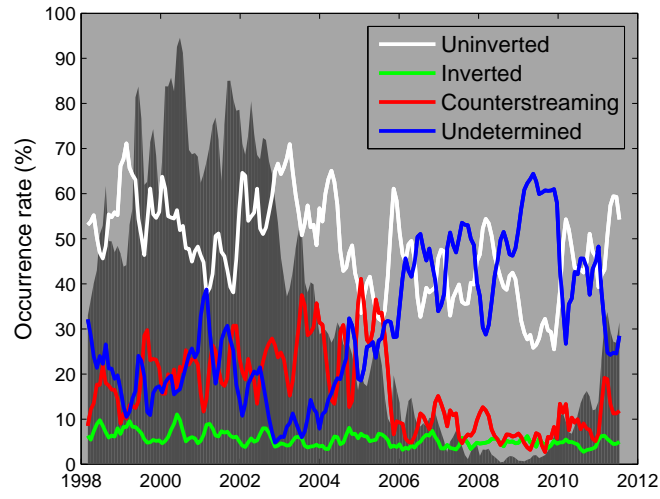


**Figure 1.** Sketches of possible HMF configurations and the resulting magnetic field and suprathermal electron signatures in near-Earth space. Red (black) arrows show the suprathermal electron strahl (magnetic field polarity), while green crosses show the position of magnetic reconnection. (a) A typical sector boundary/HCS crossing. (b) A sector boundary accompanied by closed HMF loops, likely part of an ICME. (c) A sector boundary/HCS crossing containing an inverted HMF interval at time 2. (d) An inverted HMF interval at time 2 embedded within a unipolar region. (e) A sector boundary with mismatched electron and magnetic signatures. The dashed lines show portions of the inverted HMF structure which are out of the ecliptic plane and not encountered by the observing spacecraft [after *Crooker et al.*, 2004b].

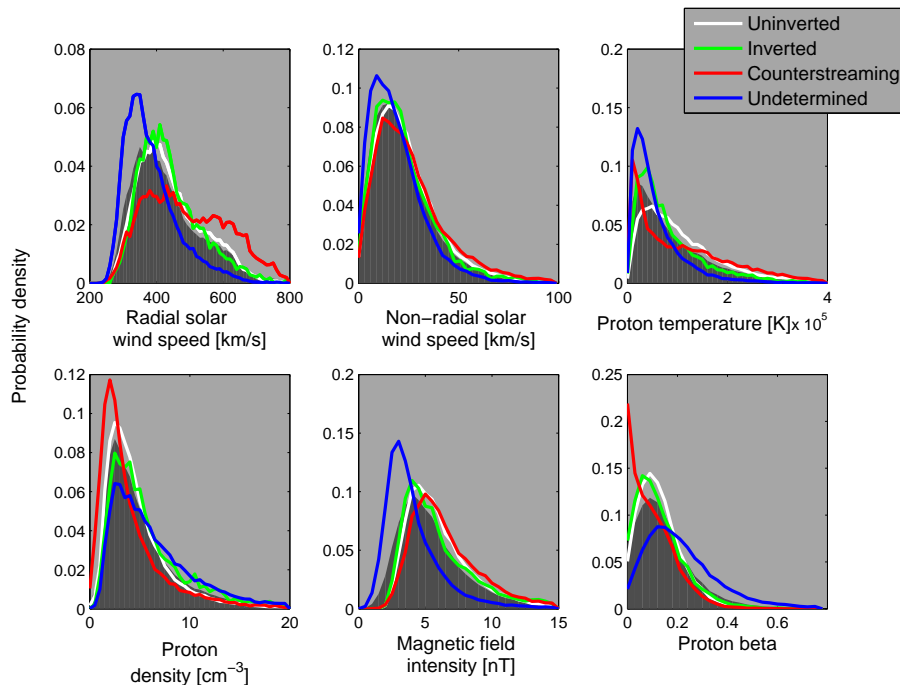
		# 1-hour intervals	% of available data
Magnetometer	Sunward HMF	53714	44.9%
Data	Antisunward HMF	56684	47.3%
	Undetermined	9366	7.83%
	Inward sector	60252	50.4%
	Outward sector	59041	49.3%
	Undetermined	371	0.31%
Suprathermal	Parallel strahl	37961	31.7%
electron data	Antiparallel strahl	37774	31.6%
	Counterstreaming	17023	14.2%
	Undetermined	26906	22.5%
Combined	Uninverted	57345	48.0%
datasets	Inverted	6608	5.53%
	Counterstreaming	19388	16.2%
	Undetermined	36139	30.2%

**Table 1.** The number of 1-hour observation periods of different HMF populations obtained using the magnetic field and suprathermal electron selection criteria.

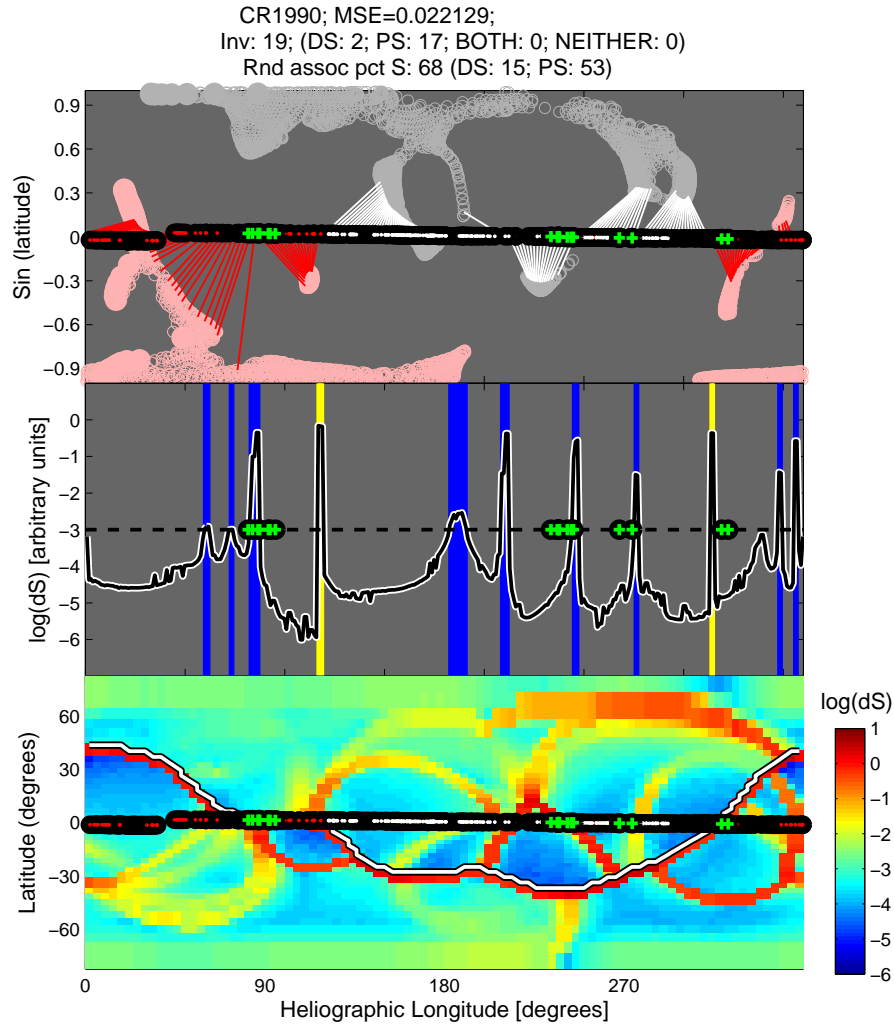




**Figure 2.** Three-Carrington rotation averages of the occurrence rates of various HMF topologies as a function of time. Sunspot number, scaled to fit the axis, is shown as the dark shaded region. Although some changes in the various HMF populations are likely to be due to changes in the electron detector, what this figure makes clear is that inverted flux is detected throughout the solar cycle.



**Figure 3.** Probability distribution functions for various near-Earth solar wind populations. The grey shaded region shows all solar wind in the interval 1998-2011. Coloured lines show subsets of these data: White, green, red and blue lines show uninverted, inverted, counterstreaming and undetermined HMF intervals, respectively.



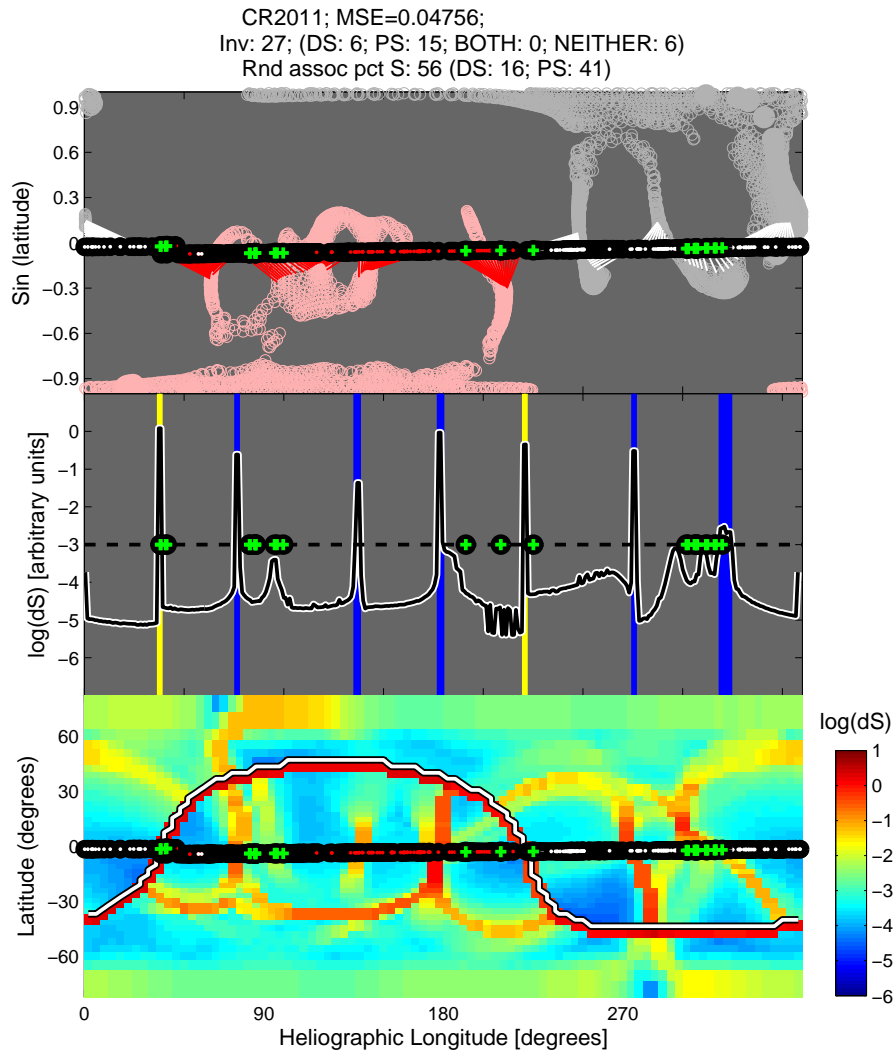
**Figure 4.** Top: A latitude-longitude map of the PFSS solution for Carrington rotation 1990. Pink/dark grey regions are the PFSS inward/outward coronal holes, with red/white lines showing the connection between the Earth's orbit across the source surface and photosphere. Overlaid on the black strip are red/white dots showing the observed outward/inward sectors mapped to the source surface. Green crosses are inverted flux intervals. Middle:  $dS$ , photospheric foot point separation for adjacent points on the source surface, along the ecliptic plane (shown on a  $\log_e$  scale). This parameter serves as a means of identifying coronal streamers: Bipolar (pseudo) streamers are shown as vertical yellow (blue) lines. Bottom: contour plot of  $dS$  over all latitudes of the source surface. The HCS

is the white curve.

D R A F T

February 7, 2013, 2:23pm

D R A F T

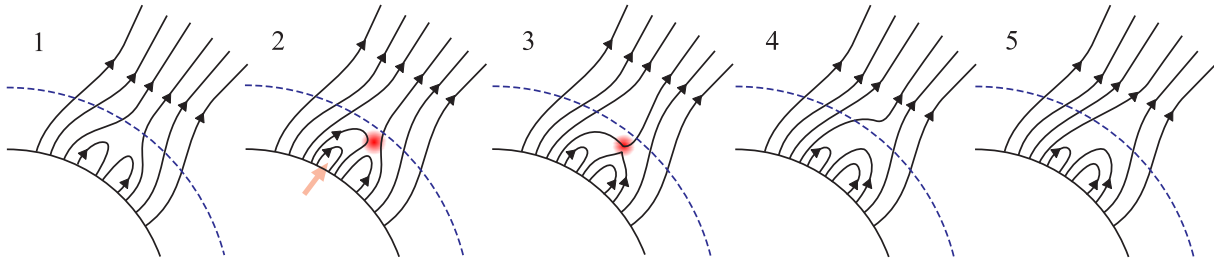


**Figure 5.** Parameters for Carrington rotation 2011, in the same format as Figure 4.

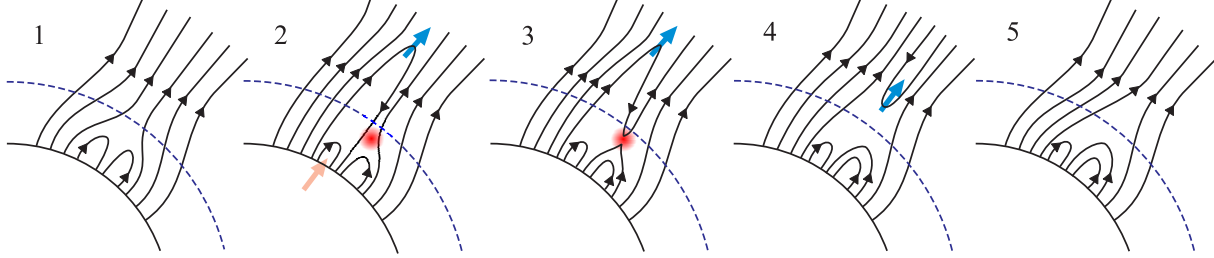
	Total	Any streamer	Pseudo (PS)	Bipolar (DS)	Both PS and DS	No streamer association
Inverted HMF	2263	1310	949	504	143	953
(% of total)	-	(57.9%)	(41.9%)	(22.3%)	(6.3%)	(42.1%)
Random interval	-	52.4%	39.0%	20.5%	5.1%	47.6%

**Table 2.** Solar origins of the inverted HMF intervals. Also shown is the probability that a random solar wind interval would be associated with the given type of streamer, *i.e.*, the percentage of ecliptic longitudes which are associated with different coronal structures.

(a) Pseudostreamer loop does not extend to solar wind formation height



(b) Pseudostreamer loop extends to solar wind formation height



**Figure 6.** A sketch of interchange reconnection within a pseudostreamer. In the top panel, a closed loop rises due to photospheric flux emergence (red arrow), but does not reach the solar wind acceleration height (blue dashed line) before it undergoes reconnection with an open magnetic field line. This creates an Alfvén wave on the open magnetic field line which propagates out into the heliosphere, but does not create inverted HMF. The bottom panels show a loop which is dragged out by the solar wind (blue arrow) before interchange reconnection occurs. It does result in the creation of inverted HMF.